

Seminar on Brown Measure (1)

25 March 2021

§1 Brown Measure for Matrices

$$T \in M_N(\mathbb{C}) \quad \text{Tr}(T) = \sum_{i=1}^N t_{ii}$$

$\lambda_1, \dots, \lambda_N$ eigenvalues of T

$$|T| = \sqrt{T^* T} \quad g(x, y) = \log \sqrt{x^2 + y^2}$$

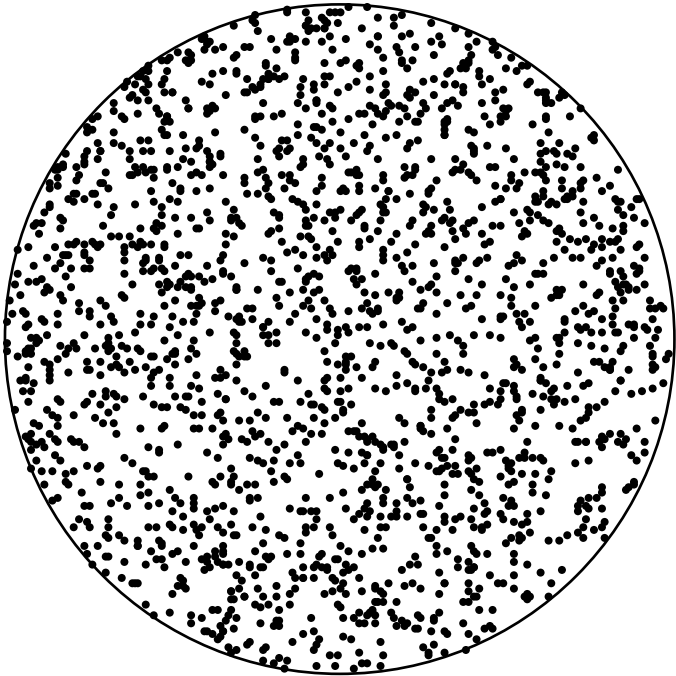
loc. integrable even on nbhd of $(0, 0)$

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \quad \text{Thm: } f \in C_c^\infty(\mathbb{R}^2)$$

$$\int \nabla^2 f(x, y) \log[\det|x+iy-T|] dx dy$$

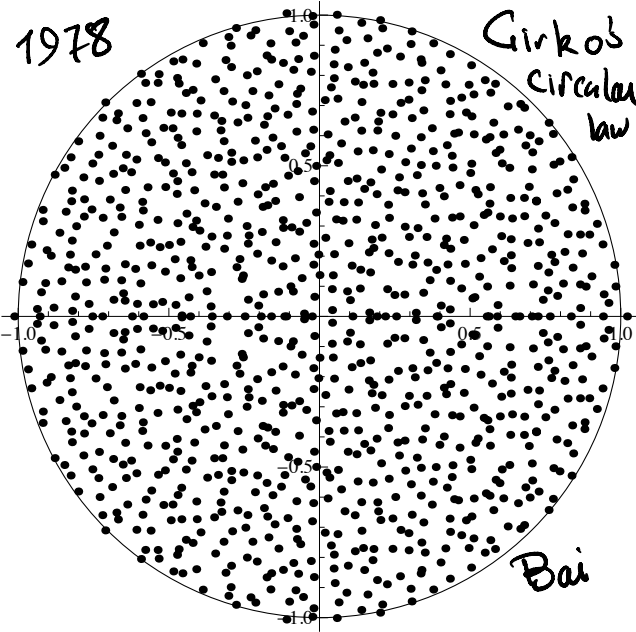
$$= 2\pi \sum_{i=1}^N f(\lambda_i) \Big|_{\lambda=x+iy}$$

$$\text{let } \mu_T = \frac{1}{2\pi N} \nabla^2 \log[\det|\lambda-T|] \\ \text{in the sense of dist.}$$

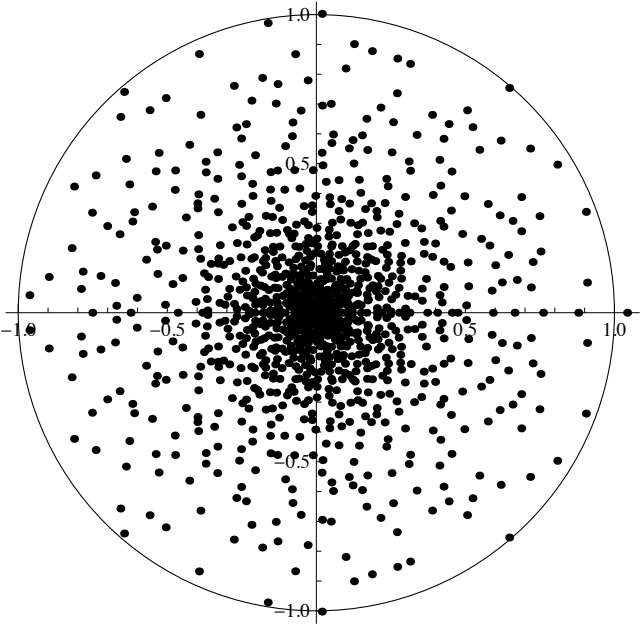


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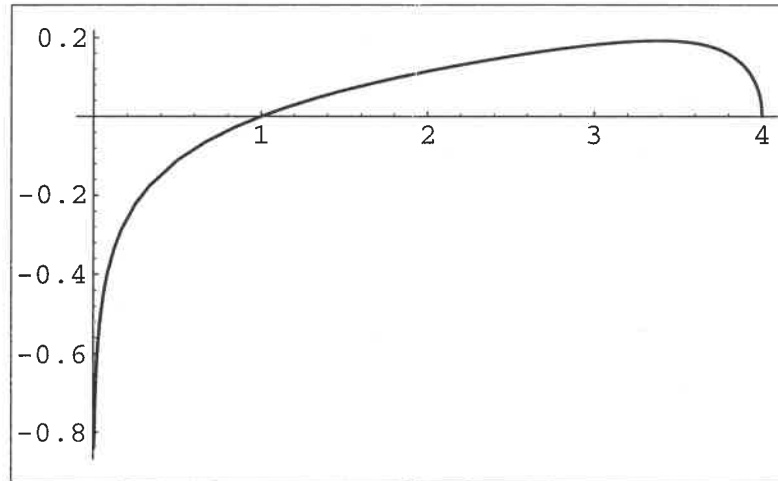
Girkos
Circular
law



Bai



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$$f(x) = \log x \frac{\sqrt{4(n-1)-x^2}}{n^2-x^2} \quad (n=5)$$

Status Report

for

Flemming Larsen

Chapter 11

Brown Measure

The Brown measure is a generalization of the eigenvalue distribution for a general (not necessarily normal) operator in a finite von Neumann algebra (i.e. a von Neumann algebra which possesses a trace). It was introduced by Larry Brown in [46], but fell into obscurity soon after. It was revived by Haagerup and Larsen [85], and played an important role in Haagerup's investigations around the invariant subspace problem [87]. By using a "hermitization" idea one can actually calculate the Brown measure by $M_2(\mathbb{C})$ -valued free probability tools. This leads to an extension of the algorithm from the last chapter to the calculation of arbitrary polynomials in free variables. For generic non-self-adjoint random matrix models their asymptotic complex eigenvalue distribution is expected to converge to the Brown measure of the ($*$ -distribution) limit operator. However, because the Brown measure is not continuous with respect to convergence in $*$ -moments this is an open problem in the general case.

11.1 Brown measure for normal operators

Let (M, τ) be a W^* -probability space and consider an operator $a \in M$. The relevant information about a is contained in its $*$ -distribution which is by definition the collection of all $*$ -moments of a with respect to τ . In the case of self-adjoint or normal a we can identify this distribution with an analytic object, a probability measure μ_a on the spectrum of a . Let us first recall these facts.

If $a = a^*$ is self-adjoint, there exists a uniquely determined probability measure μ_a on \mathbb{R} such that for all $n \in \mathbb{N}$

$$\tau(a^n) = \int_{\mathbb{R}} t^n d\mu_a(t)$$

and the support of μ_a is the spectrum of a ; see also the discussion after equation (2.2) in Chapter 2.

More general, if $a \in M$ is *normal* (i.e., $aa^* = a^*a$), then the spectral theorem provides us with a projection valued spectral measure E_a and the Brown measure is

just the spectral measure $\mu_a = \tau \circ E_a$. Note that in the normal case μ_a may not be determined by the moments of a . Indeed, if $a = u$ is a Haar unitary then the moments of u are the same as the moments of the zero operator. Of course, their $*$ -moments are different. For a normal operator a its spectral measure μ_a is uniquely determined by

$$\tau(a^n a^{*m}) = \int_{\mathbb{C}} z^n \bar{z}^m d\mu_a(z) \quad (11.1)$$

for all $m, n \in \mathbb{N}$. The support of μ_a is again the spectrum of a .

We will now try to assign to any operator $a \in M$ a probability measure μ_a on its spectrum, which contains relevant information about the $*$ -distribution of a . This μ_a will be called the *Brown measure* of a . One should note that for non-normal operators there are many more $*$ -moments of a than those appearing in (11.1). There is no possibility to capture all the $*$ -moments of a by the $*$ -moments of a probability measure. Hence, we will necessarily lose some information about the $*$ -distribution of a when we go over to the Brown measure of a . It will also turn out that we need our state τ to be a trace in order to define μ_a . Hence in the following we will only work in tracial W^* -probability spaces (M, τ) . Recall that this means that τ is a faithful and normal trace. Von Neumann algebras which admit such faithful and normal traces are usually addressed as *finite* von Neumann algebras. If M is a finite factor, then a tracial state $\tau : M \rightarrow \mathbb{C}$ is unique on M and is automatically normal and faithful.

11.2 Brown measure for matrices

In the finite-dimensional case $M = M_n(\mathbb{C})$, the Brown measure μ_T for a normal matrix $T \in M_n(\mathbb{C})$, determined by (11.1), really is the eigenvalue distribution of the matrix. It is clear that in the case of matrices we can extend this definition to the general, non-normal case. For a general matrix $T \in M_n(\mathbb{C})$, the spectrum $\sigma(T)$ is given by the roots of the characteristic polynomial

$$P(\lambda) = \det(\lambda I - T) = (\lambda - \lambda_1) \cdots (\lambda - \lambda_n),$$

where $\lambda_1, \dots, \lambda_n$ are the roots repeated according to algebraic multiplicity. In this case we have as eigenvalue distribution (and thus as Brown measure)

$$\mu_T = \frac{1}{n}(\delta_{\lambda_1} + \cdots + \delta_{\lambda_n}).$$

We want to extend this definition of μ_T to an infinite dimensional situation. Since the characteristic polynomial does not make sense in such a situation we have to find an analytic way of determining the roots of $P(\lambda)$ which survives also in an infinite dimensional setting.

Consider

$$\log |P(\lambda)| = \log |\det(\lambda I - T)| = \sum_{i=1}^n \log |\lambda - \lambda_i|.$$

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FASCICULE XXV

LES ALGÈBRES D'OPÉRATEURS
DANS L'ESPACE HILBERTIEN
(ALGÈBRES DE VON NEUMANN)

PAR

Jacques DIXMIER

PROFESSEUR A LA FACULTÉ DES SCIENCES DE PARIS

DEUXIÈME ÉDITION
revue et augmentée

(1st ed 1957)

gv

PARIS
GAUTHIER-VILLARS ÉDITEUR
1969

$$\mu_T = \frac{1}{2\pi N} \nabla^2 \log |\det [x+iy-T]| \quad (2)$$

$$= \frac{1}{2\pi} \nabla^2 \operatorname{tr}(\log |x+iy-T|)$$

$$\operatorname{tr}(T) = \frac{1}{N} \operatorname{Tr}(T)$$

§2 Extensions to von Neumann

H = sep. Hilbert space, $B(H)$

= bdd. linear ops. $H \rightarrow H$.

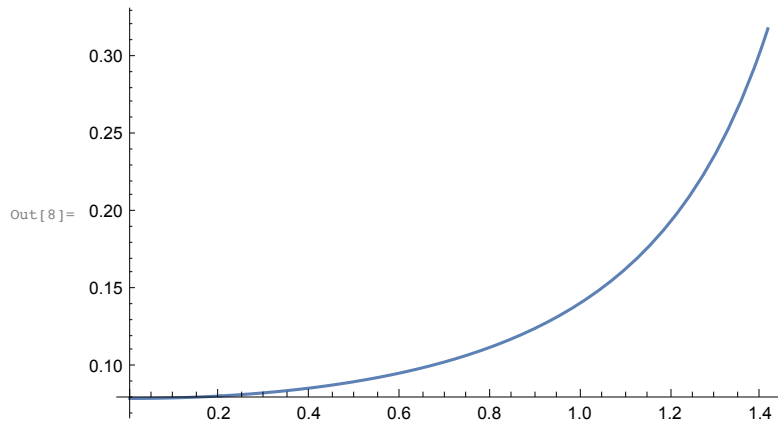
$$\xi, \eta \in H \quad \{T_n\}_n \subseteq B(H)$$

$$T_n \xrightarrow{\text{WOT}} T \quad \text{if } \forall \xi, \eta$$

$$\langle (T_n - T)\xi, \eta \rangle \rightarrow 0$$

$M \subseteq B(H)$ self-adjoint, contains \mathbb{I} , WOT
closed sub-algebra = von Neumann
alg.

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In[8]:= Plot[ $\frac{4}{\pi (4 - r^2)^2}$ , {r, 0,  $\sqrt{2}$ }]
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In[10]:= 2  $\pi$  Integrate[r  $\frac{4}{\pi (4 - r^2)^2}$ , {r, 0,  $\sqrt{2}$ }]
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Out[10]= 1

$M \subseteq B(H)$ v.N. alg.

$\tau: M \rightarrow \mathbb{C}$, linear

• $\tau(ab) = \tau(ba) \quad \forall a, b \in M$

• $\tau(x^*x) \geq 0 \quad \forall x \in M$

• $\tau(x^*x) = 0 \Leftrightarrow x = 0$

• $\tau(1) = 1$

• WOT cont.: $x_n \xrightarrow{\text{WOT}} x$

$\Rightarrow \tau(x_n) \rightarrow \tau(x)$

$(M_N(\mathbb{C}), \tau)$, $(L^\infty(\mu), S)$

$M_2 \xrightarrow{\quad} M_2 \otimes M_2 \quad x \mapsto \begin{pmatrix} x & 0 \\ 0 & x \end{pmatrix}$

$\bigotimes_{n=1}^{\infty} M_2 = M_{2^{\infty}}$

§ 3 Examples

s_1, s_2 free & semi-circular

$$c = \frac{1}{\sqrt{2}} (s_1 + i s_2) \quad \text{Voiculescu's circular op.}$$

$\mu_c = \text{uniform measure on } \mathbb{T}$

Proof of Thm

$f \in C_c^\infty(\mathbb{R}^2)$ claim:

$$\int \nabla^2 f(x, y) \log \sqrt{x^2 + y^2} \, dx \, dy$$

$$= 2\pi f(0, 0)$$

Choose $R > 0$ $\text{supp}(f) \subseteq D_R$ open

$$D_{r, R} = \{ (x, y) \mid r < \sqrt{x^2 + y^2} < R \}$$

Calculus Lemma: $\nabla \cdot (f \nabla g - g \nabla f)$
 $= f \nabla^2 g - g \nabla^2 f$

$$g(x, y) = \log \sqrt{x^2 + y^2} \quad \nabla^2 g = 0$$

$$\nabla g = \left(\frac{x}{x^2 + y^2}, \frac{y}{x^2 + y^2} \right) \quad \text{on } \partial D_r$$

let \vec{n} = outward pointing normal

$$= \left(\frac{x}{\sqrt{x^2 + y^2}}, \frac{y}{\sqrt{x^2 + y^2}} \right)$$

$$\nabla g \cdot \vec{n} = \frac{1}{r} \quad \text{on } \partial D_r$$

$$\int_{D_{r,R}} \nabla^2 f g \, dx \, dy = \int_{D_{r,R}} \nabla \cdot (g \nabla f - f \nabla g) \, dx \, dy$$

$$= \int_{\partial D_{r,R}} (g \nabla f - f \nabla g) \cdot \vec{n} \, ds$$

$$= \int_{\partial D_{r,R}} g \nabla f \cdot \vec{n} ds - \int_{\partial D_{r,R}} f \nabla g \cdot \vec{n} ds$$

$$= \frac{1}{r} \int_0^{2\pi} r^2 f(r \cos \theta, r \sin \theta) d\theta$$

$$- r \log r \int_0^{2\pi} \frac{\partial f}{\partial r}(r \cos \theta, r \sin \theta) d\theta$$

$$\rightarrow 2\pi f(0,0) \quad \square$$

$$\det(\lambda - T) = \prod_{i=1}^N (\lambda - \lambda_i)$$

$$\det |\lambda - T| = |\det(\lambda - T)|$$

$$= \prod_{i=1}^N |\lambda - \lambda_i|$$

$$\log(\det |\lambda - T|) = \sum_{i=1}^N \log |\lambda - \lambda_i|$$

||

$$\operatorname{Tr}(\log |\lambda - T|)$$

$$\frac{1}{2\pi} \operatorname{tr}(\log |\lambda - T|) = \frac{1}{2\pi N} \sum_{i=1}^N \log |\lambda - \lambda_i|$$

$$\int \nabla^2 f(x, y) \left[\frac{1}{2\pi} \operatorname{tr}(\log |\lambda - T|) \right] dx dy$$

$$= \frac{1}{2\pi N} \sum_{i=1}^N \int \nabla^2 f(x, y) \log |\lambda - \lambda_i| dx dy$$

$$= \frac{1}{N} \sum_{i=1}^N f(\lambda_i)$$

$$= \int f(\lambda) d\mu_T(\lambda)$$

$$d\mu_T = \frac{1}{N} \sum_{i=1}^N \delta_{\lambda_i}$$

V is a vector space

$P: V \rightarrow V$ is a linear transformation such

that $P^2 = P$ i.e. $P(P(v))$

$= P(v)$, P is a projection.

$$\ker(P) = \{v \mid P(v) = 0\}$$

$$\text{ran}(P) = \{P(v) \mid v \in V\}$$

given $v \in V \exists w, z \in V$

$w \in \text{ran}(P)$, $z \in \ker(P)$, and

$$v = w + z.$$

$$w = P(v) \quad z = (I - P)(v) = v - w$$

$$w + z = P(v) + (I - P)(v)$$

$$= (P + I - P)(v) = I(v) = v$$

$$P(z) = P(I - P)(v) = (P - P^2)(v)$$

$$= 0 \Rightarrow z \in \ker(P)$$

$$P(w) = P(P(v)) = P(v) = w$$

$$\uparrow \\ \text{ran}(P)$$

